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PROBABILITY DISTRIBUTION
OF IMPACT FOR A SATELLITE
BOOSTER AFTER LOSS OF CONTROL

THESIS

AFIT/GA/AA/78D-1 Scott W. Benson 2nd Lt . USAF



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THESTS

of the Air Force Institute of Technology

Air Training Command

in Partial Fulfillment of the

Requirements for the Degree of

Master of Science

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USAF

Graduate Astronautical Engineering

1.144

March 1979

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Acknowledgement

I wish to thank my wife, Pam, for her understanding and her persistence in keeping me at work. I also wish to thank my advisor, Capt. Bill Wiesel, for his advice and patience towards the completion of this study.

Scott W. Benson

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List of Symbols

Symbol	Definition
а	Semi-major axis of conic section
е	Eccentricity of conic section
E	Probability density
E,	Eccentric anomaly at impact
E _o	Eccentric anomaly at failure
g	Gravitational acceleration
h	Altitude
ħ	Angular momentum
m	Mass
p	Parameter of conic section
P	Probability
ř	Position of vehicle
R _{PER}	Ferigee height of conic section
$\mathtt{R}_{\mathbf{E}}$	Radius of earth
t	Time
$\mathtt{t}_\mathtt{B}$	Rocket burn time
T	Thrust
\vec{v}	Velocity of vehicle
V _{CIRC}	Velocity to obtain circular orbit
${f v}_{f E}$	Equivalent exit velocity of rocket
v _R ·	Residual velocity
x	Range
α	Azimuth
β	Mass flow rate
Y	Flight rath inclination

Symbol	Definition
δ	Elevation
0	Longitude
å	Angle of inclination
μ	Earth gravitational constant
v ₆	True anomaly at impact
o ^v o	True anomaly at failure
φ	Latitude
J	Angle traversed by vehicle

Abstract

The impact area and regions of high impact probability after booster failure were found for a sample space system. The impact footprints were found to be drop-shaped. They increased in size as the residual velocity vector elevation approached the original flight path, or when the booster failure was assumed earlier in the boost trajectory. Probability densities at impact were found, and plotted against the longitude of impact. Singularities in the probability density were discovered. The singularities correspond to areas of high impact probability, and occur when a maximum latitude or longitude is obtained at a fixed residual velocity vector azimuth or elevation. These singularities were plotted along the impact envelopes. A sample failing booster system was used to demonstrate the impact area and probability characteristics.

PROBABILITY DISTRIBUTION OF IMPACT FOR A SATELLITE BOOSTER AFTER LOSS OF CONTROL

I. Introduction

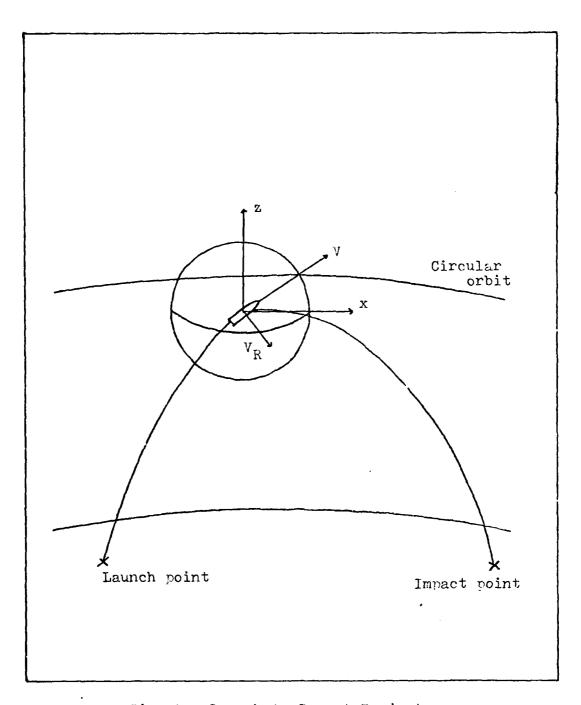
Background

In light of the recent Soviet satellite reentry and impact in Canada, it is becoming more important to determine the impact areas and probabilities of failing space systems. The main reasons for concern are the possibilities that there are nuclear materials on board, in which personal safety is endangered or that the system is highly classified, in which national security is endangered. Of special interest to the United States are the slightly retrograde orbits launched from Vandenburg AFB, California. These type of orbits pass over Moscow in the first revolution, and might be cause for concern in the case of booster failure.

The objective of this study was to demonstrate a method of predicting impact areas and probabilities for a failing booster system. A booster is launched and, at some time before orbit injection, suffers total loss of control. The vehicle follows a new trajectory to earth impact. The position of impact and the probability is calculated for each possible failure. The launch to impact trajectory is represented in Fig 1.

Approach

The booster equations of motion were first integrated to a circular orbit. This was accomplished by setting final



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Fig. 1. Launch-to-Impact Trajectory

orbit conditions and solving the boundary value problem. Once a launch-to-orbit trajectory was obtained, the equations of motion were integrated to several times lower than the time to orbit injection. At this point, the residual velocity was applied in one of many failure directions. The vehicle position and new vehicle velocity were obtained and the orbit they define was found. From the orbital elements, it was determined when and where the vehicle impacted the earth. Knowing the impact point, the probability density at impact was found by propagating the probability density at failure, through the new trajectory, to the ground.

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Scope

For this initial study, the scope of the problem was limited. The booster was assumed to launch easterly along the equator, through a vacuum, to an arbitrary circular orbit height. The booster was modeled as a simple single stage booster, with a constant thrust profile. The number of possible failures was limited to a small, representative sample. There are several other less important assumptions made throughout the body of this report.

The following chapter addresses the problem solving procedure in detail.

II. Procedure

Launch to Orbit Injection

The beginning stategy was to find a launch trajectory that terminated in a circular orbit. Obtaining this trajectory requires the solution of a boundary value problem, using the equations of motion for a gravity turn trajectory (Ref 5:335-43).

$$\dot{x} - V\cos\gamma = 0$$

$$\dot{h} - V\sin\gamma = 0$$

$$\dot{V} + g\sin\gamma - \frac{\beta V_E}{m} = 0$$

$$\dot{\gamma} + \frac{g\cos\gamma}{V} = 0$$

$$\dot{n} + \beta = 0$$

$$T = \beta V_E$$

In these equations, x and h are the vehicle range and altitude respectively, V is the velocity magnitude, γ is the flight path inclination with respect to the local horizontal, m is the vehicle mass, β is the rocket mass flow rate, and V_E is the rocket equivalent exit velocity. The launch phase is illustrated in Fig 2.

There are several underlying assumptions used in obtaining this set of equations. First, the trajectory is relative to a normal reference frame; therefore, the local earth must be assumed to be flat. Second, flight is assumed to be in a vacuum, negating the lift and drag effects. Thirdly, the thrust is assumed to act parallel to the velocity.

Other assumptions were made to simplify the boundary

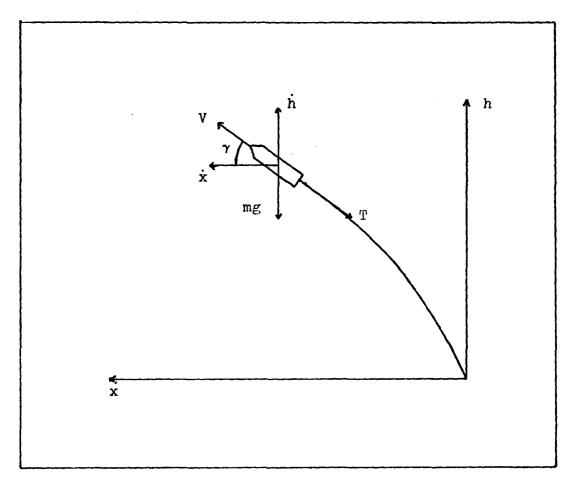


Fig. 2. Booster Launch

value launch problem. The mass flow rate, β , and the equivalent exit velocity, V_E , were assumed to be constant, resulting in a constant thrust profile. The gravitational acceleration, g, was also assumed to be constant.

The boundary value launch was executed as follows. An initial mass, and appropriate values of thrust, mass flow rate, and equivalent exit velocity were chosen. The equations of motion were integrated through a number of different burn times, and the resulting flight path inclinations

were obtained. The integration routine used was ODE, an ordinary differential equation integrator found in the IMSL subroutine package. The burn time to zero path inclination, when the booster is parallel to the earth, was found using the Aitken method of interpolation. The Aitken subroutine, INTERP, is also found in the IMSL library. Integrating again to this burn time gives values for path inclination, velocity and altitude. To achieve successful circular orbit, the final velocity must be equal to the circular orbit velocity at that altitude. The initial value of mass was changed to obtain this required boundary condition.

The method used in changing the initial condition, mass, is discussed here. Rearranging the equations of motion gives an equation for instantaneous change in velocity for a change in mass.

$$\frac{dV}{dm} = \frac{-T}{\beta m} + g \frac{\sin \gamma}{\beta}$$

The required change in velocity is

$$\Delta V = V - V_{CIRC}$$

where

$$V_{CIRC} = \sqrt{\frac{\mu}{R_E + h}}$$

The first change in mass becomes

$$\Delta m = \frac{-\Delta V}{\frac{dV}{dm}}$$

This change in mass was added a number of times; the equations of motion being integrated to a horizontal flight path for each mass increment. Again, the Aitken intervola-

tion scheme was used to find the initial value of the total mass which satisfied the boundary conditions.

The interpolated values of mass and burn time were then used to integrate the equations of motion to the final set of parameters.

Booster Failure

Booster failure was modeled by an instantaneous change in direction of the residual velocity at a given time in the booster trajectory.

A time decrement was chosen and subtracted from the circular orbit trajectory burn time. The equations of motion were integrated to the resultant time, which gave a set of parameters before orbit injection. The magnitude of the residual velocity was found by subtracting the final velocity Verre at with of from the circular orbit velocity.

$$v_R = v_{CIRC} - v$$

It was assumed that the direction of the residual velocity could be towards any point on a sphere surrounding the vehicle. For this study, the direction was incremented by 30° in azimuth and elevation, to give a total of 62 possible directions. The increment was decreased later in the study to achieve a better understanding of the probabilities of impact. The sphere of points was defined with the polar axis along the local vertical, as seen in Fig 3.

The components of the residual velocity can easily be found, as shown in Fig 4.

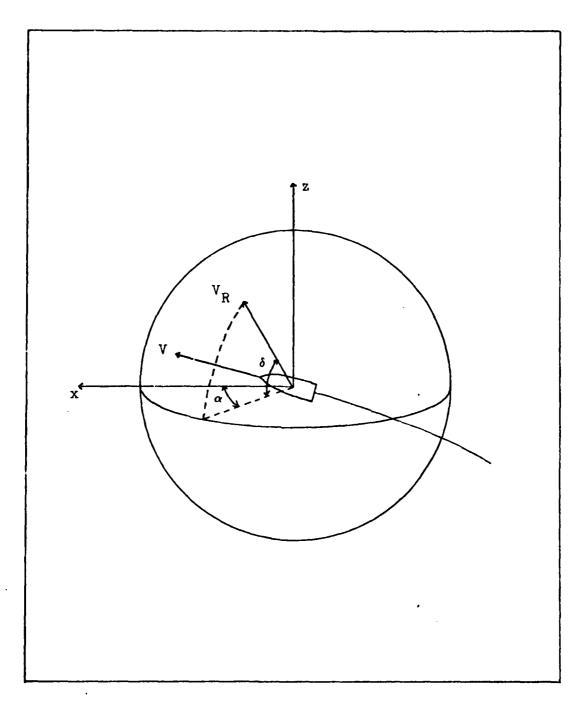


Fig. 3. Booster Failure

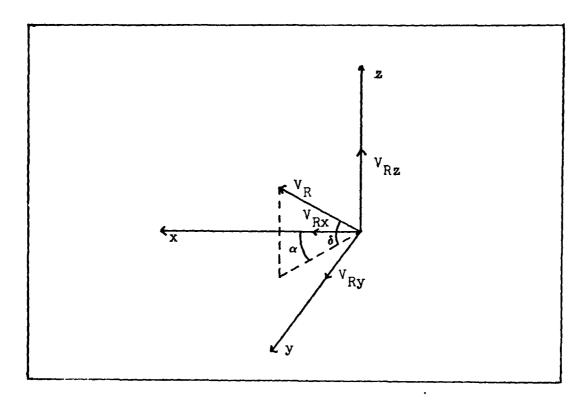


Fig. 4. Residual Velocity Components

 $V_{Rx} = V_{R} \cos \alpha \cos \delta$

 $V_{Ry} = V_{R} \sin \alpha \cos \delta$

 $V_{Rz} = V_{R} \sin \delta$

where

the z axis is aligned along the local vertical the x axis is in the plane of the original flight path

These components were then added to the flight velocity components, \dot{x} and \dot{h} , at burnout to give a resultant velocity after booster failure, herein referred to simply as V. This velocity and the vehicle position, as found through the equation integration, gives the information needed to propagate the trajectory to impact.

Flight Trajectory Propagation to Impact

The requirements for this part of the study were to see if the booster impacts the earth, and if so, where it impacts and how long a time interval it covers before impact.

The method used in finding these answers starts with the conversion of the vehicle velocity and position after failure to the orbital elements. Using the orbital elements, it is simple to check if the vehicle impacts by checking the orbit perigee height. If it does impact, the angle traversed and the angle of deviation from the original flight path can be found. Knowing the angle traversed, the time of flight can be calculated, and the impact latitude and longitude are easily obtained.

The vehicle velocity and position after failure were converted to the necessary orbital elements (Ref 3:17-26).

$$\vec{h} = \vec{r} \times \vec{V}$$

$$\vec{e} = \frac{1}{\mu} \left[\left(V^2 - \frac{\mu}{r} \right) \vec{r} - \left(\vec{r} \cdot \vec{V} \right) \vec{V} \right]$$

$$p = \frac{h^2}{\mu}$$

$$a = \frac{p}{1 - e^2}$$

Using the calculated parameter and eccentricity, the orbit perigee height was found and checked against the earth radius, which is assumed to be constant.

$$R_{\text{PER}} = \frac{p}{1 + e}$$

check

$$R_{FER} \leq R_{E}$$

If the perigee height was greater than the earth radius, the

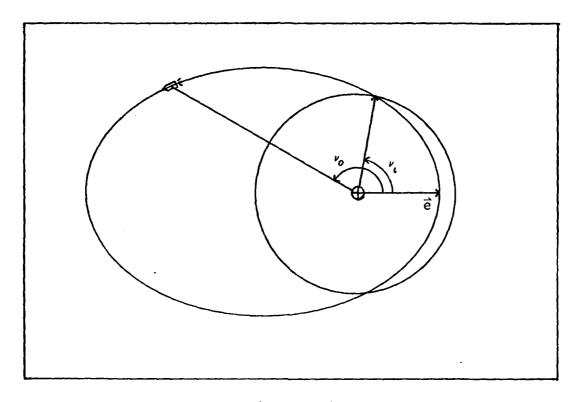


Fig. 5. Vehicle Orbit to Impact vehicle would enter earth orbit, and those cases were ignored.

If the perigee check indicated that impact did occur, the angle the vehicle traverses before impact, ψ , can be found by

$$\psi = 2\pi - \nu_0 - \nu_t$$

where

", is the true anomaly at failure
", is the true anomaly at impact, as seen in Fig 5

These true anomalies can be calculated using the equation of a conic section.

$$\nu_{o} = \cos^{-1}\left(\frac{\frac{p}{r}-1}{e}\right)$$

$$\nu_{i} = \cos^{-1}\left(\frac{\frac{p}{R_{E}}-1}{e}\right)$$

The angle of inclination, , after failure is found by

$$\iota = \cos^{-1} \frac{\vec{h} \cdot \hat{k}}{\vec{h}}$$

where \hat{k} is the unit vector along the z axis.

To obtain the time of flight from failure to impact, the eccentric anomaly, E, at each point of interest must be found.

$$E_{o} = \cos^{-1} \left(\frac{e + \cos \nu_{o}}{1 + e \cos \nu_{o}} \right)$$

$$E_{t} = \cos^{-1} \left(\frac{e + \cos(2\pi - \nu_{t})}{1 + e \cos(2\pi - \nu_{t})} \right)$$

The time of flight is found by solving Kepler's equation for each eccentric anomaly, and getting the difference in time.

$$\Delta t = \sqrt{\frac{3}{\mu}} \left[(E_i - esinE_i) - (E_o - esinE_o) \right]$$

This change in time represents the time from failure to impact. This time is added to the burn time to failure to obtain the total time of flight.

Using spherical trigonometry, the latitude, ϕ , and longitude, θ , are easily found from ψ and ι , as shown in Fig 6. The original simplifying assumption of launching along the equator takes effect here.

$$\phi = \sin^{-1}(\sin_i \sin \psi)$$

$$\theta = \sin^{-1}\frac{\tan \phi}{\tan \psi}$$

In the cases when the change in velocity occurs in the plane of the original orbit, the latitude and longitude are simply

$$\phi = 0$$

$$\theta = \psi$$

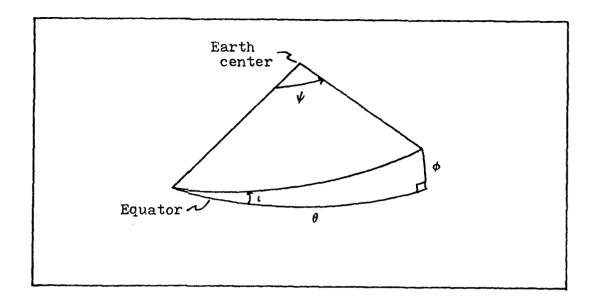


Fig. 6. Latitude and Longitude Traversed The longitude can be added to the longitude covered prior to failure, $\theta_{\rm f}$, to obtain a total longitude. The fail-

ure longitude is found with the original flat earth assump-

tion.

$$\theta_f = \sin^{-1} \frac{x}{R_E}$$

$$\theta_t = \theta + \theta_f$$

At this point, a position and time of impact has been found for a vehicle which had been subjected to a velocity change before injection. The probability of impact at this position will now be determined.

Probability Determination

To determine the probability of impact at a specific latitude and longitude, the probability density at failure for the corresponding velocity change is found, then propagated along the trajectory to the ground, as seen in Fig 7.

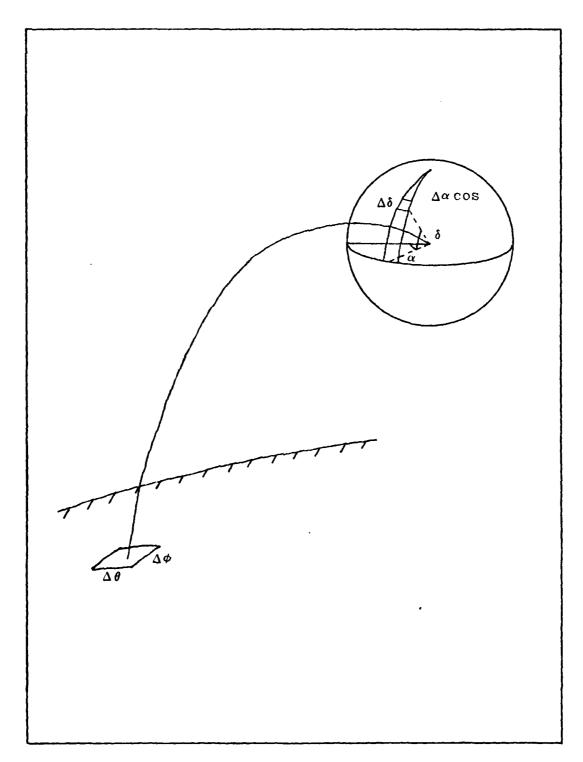


Fig. 7. Probability Propagation to Impact

For a small change in azimuth and elevation, the probability density at failure is

$$E(\alpha,\delta) = \frac{\cos\delta}{4\pi}$$

Therefore, the probability of the vehicle massing through an area of the sphere is

$$F(\alpha, \delta) = \int_{\alpha}^{\alpha + \Delta \alpha} d\alpha \int_{\delta}^{\delta + \Delta \delta} d\delta \frac{\cos \delta}{4\pi}$$

When integrated over the whole sphere (from $0 \rightarrow 2\pi$ in azimuth and $-\frac{\pi}{2} \rightarrow \frac{\pi}{2}$ in elevation), the probability becomes unity.

Fromagation of the probability to a latitude and longitude at impact results in

$$P(\phi,\theta) = \int_{\Phi}^{\phi + \Delta\phi} d\phi \int_{\theta}^{\theta + \Delta\theta} d\theta \frac{\cos\delta}{4\pi} \frac{\partial(\alpha,\delta)}{\partial(\phi,\theta)}$$

The Jacobian is defined as a determinant.

$$\frac{\partial(\alpha,\delta)}{\partial(\phi,\theta)} = \begin{vmatrix} \frac{\partial\alpha}{\partial\phi} & \frac{\partial\delta}{\partial\phi} \\ \frac{\partial\alpha}{\partial\theta} & \frac{\partial\delta}{\partial\theta} \end{vmatrix}$$

As the changes in latitude and longitude become infinitesmal, the probability density becomes the integrand.

$$E(\phi,\theta) = \frac{\cos\delta}{4\pi} \begin{vmatrix} \frac{\partial\alpha}{\partial\phi} & \frac{\partial\delta}{\partial\phi} \\ \frac{\partial\alpha}{\partial\theta} & \frac{\partial\delta}{\partial\theta} \end{vmatrix}.$$

In this study, the partial derivatives are estimated by numerical methods. Very small changes in azimuth and elevation are chosen and added. When the computer program is run, there are resulting small changes in latitude and longitude. The ratio of these changes becomes the numerical partial derivative.

$$E(\phi,\theta) = \frac{\cos \delta}{4\pi} \begin{vmatrix} \frac{\Delta \alpha}{\Delta \phi} & \frac{\Delta \delta}{\Delta \phi} \\ \frac{\Delta \alpha}{\Delta \theta} & \frac{\Delta \delta}{\Delta \theta} \end{vmatrix}$$

This probability density represents probability in relative magnitude. When the density is high, or approaching a singularity, the probability is high for that impact point.

The following chapter will discuss the resulting impact envelopes and probability density behavior for this study.

III. Results and Discussion

Gravity-turn Trajectory

To initiate the gravity-turn launch, the following parameters were chosen.

$$m = 21,000 lb_{m}/g$$

$$\beta = .466 \text{ lb}_{\text{m}}/\text{sec}$$

$$T = 58,800 \text{ lb}_f$$

The launch-to-orbit trajectory had these final values.

 $t_{\rm R} = 276.54996 \, {\rm sec}$

x = 3078159.92872 ft

h = 526417.72579 ft

V = 25616.04328 ft/sec

 $\gamma = .00000036 \text{ rad} \approx 0$

 $m = 398.14796 lb_{m}$

Impact Areas

Impact points for each possible failure in each time step were found. The plot in Fig 8 is a representative plot using one time decrement and elevations from $-\frac{\pi}{2} \to \frac{\pi}{2}$ radians in $\frac{\pi}{6}$ radian increments. Each curve of points represents a constant elevation of the residual velocity vector, with azimuth rotating through $\frac{\pi}{2}$ radians. It is seen that minimum latitude and longitude dispersions occur for elevations approaching the poles of the sphere of failure. The maximum dispersions occur for elevations equal to zero.

From the proximity of the impact points for each positive and negative elevation trajectory (i.e. ±30 deg), it is

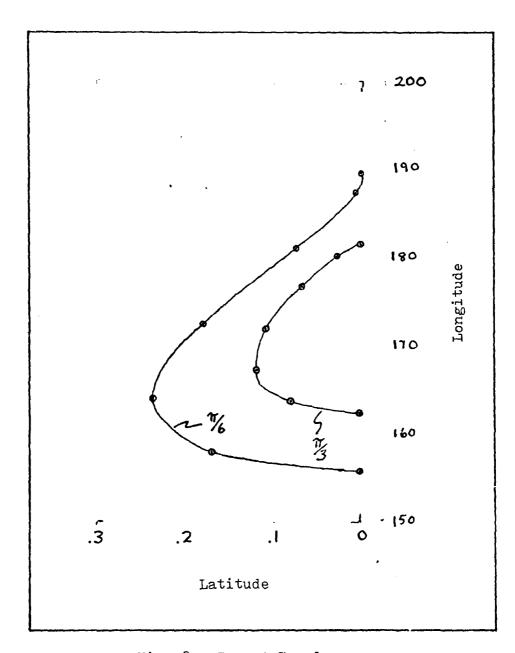


Fig. 8. Impact Envelopes

assumed that there are high and low trajectories to each impact point, excepting the case when the elevation is zero. This point is important in finding a probability for any impact point, since the probability density is composed of two probability densities summed.

In Fig 9, the maximum impact envelopes are plotted for five time decrements from orbit injection time. These represent each time step with the residual velocity vector elevation equal to zero. For the launch trajectory utilized in this study, failure results in a large dispersion in longitude and a relatively small dispersion in latitude. As the time decrements are subtracted, these dispersions grow. These characteristics would be expected in any launch configuration.

Probability Distribution

In an attempt to uncover curves of equal probability, the probability density was plotted versus the impact longitude. The result, for a specific time decrement and residual velocity vector elevation, is shown in Fig 10. A number of singularities were discovered in the probability density. These singularities represent areas of high probability.

The end singularities occur when the residual velocity vector azimuth is 0 and * radians. Close to these azimuths, a small change in residual velocity vector elevation results in a very small change in impact latitude, so there is a higher probability of impacting in those areas.

The middle singularity corresponds to an azimuth of $\frac{\pi}{2}$

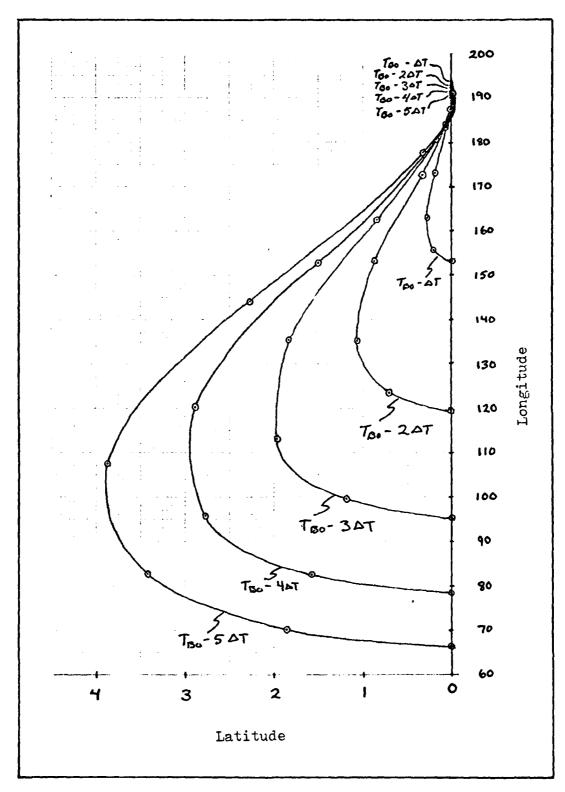


Fig. 9. Immact Envelopes with Changing Time of Failure

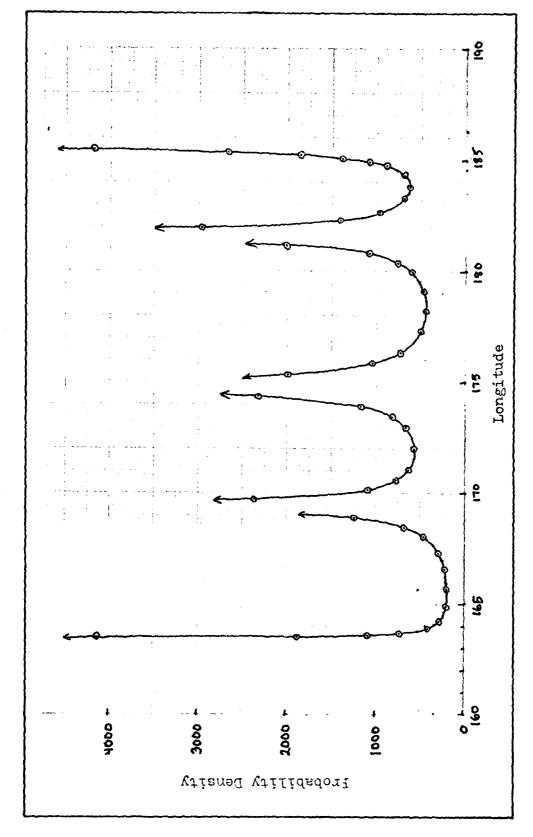


Fig. 10. Probability Density Along an Impact Envelope

radians. Two phenomena occur near this point. A small change in latitude, and a small change in azimuth result in a negligible change in longitude. These occurences give this area a high probability of impact.

The remaining singularities occur at maximum latitudes at set conditions. The singularity at ≈ 181.5 degrees corresponds to a maximum latitude for a constant azimuth of $\frac{5\pi}{18}$ radians. The singularity at ≈ 169.4 degrees corresponds to the maximum latitude for the elevation used for this plot. In these areas, a small change in elevation and azimuth, respectively, results in a negligible change in latitude, increasing the probability of impact. The position of the singularities along the impact curves, for the case in Fig 10, can be seen in Fig 11.

The number of singularities occurring can now be realized. Taking each burn time separately, for each set azimuth or elevation, there is a maximum latitude and longitude attainable, and a singularity corresponding to each. It would also be expected to have singularities corresponding to maximum latitudes and longitudes for a fixed residual velocity magnitude. All these singularities could be plotted, as in Fig 11, to obtain curves of high probability of impact for each time decrement. This is an area where further study would be enlightening, along with those mentioned in the next chapter on reccomendations.

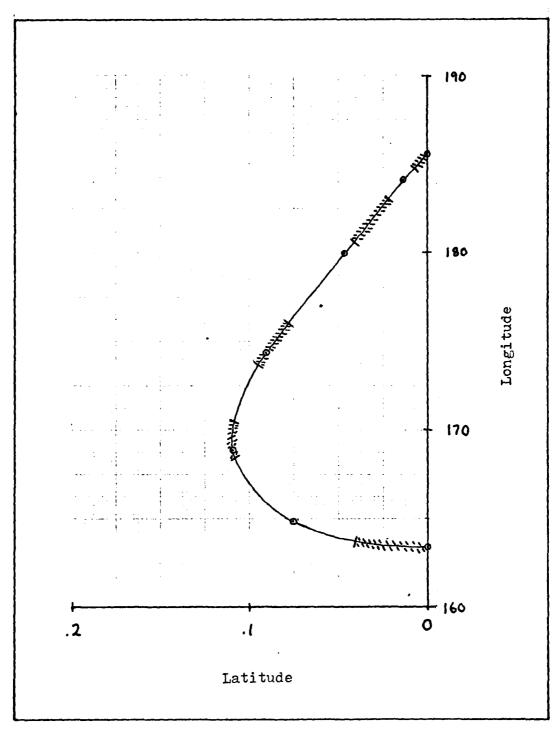


Fig. 11. Regions of High Probability Along an Impact Envelope

IV. Recommendations

There are many candidate areas for further study in this thesis. The scope of the study could be expanded to work with specific launch vehicles or launch sites. The launch equations of motion could be changed to incorporate a spherical earth, a changing gravitational acceleration, a non-constant thrust profile, a multi-stage vehicle, or flight through an atmosphere. The boundary value conditions could be changed to achieve an elliptical orbit.

As mentioned in the results chapter, the study could be taken deeper, without changing the scope, to find curves of high probability. Another approach to finding probability distributions would be to start at impact and work back to failure. This method would eliminate the problem of summing probabilities from near coincedent high and low trajectories. The depth in which the interested reader wishes to follow the study is certainly dependant on his desired results.

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Appendix A

Boundary Value Launch Computer Program

```
FROGRAM LOV(INPUT= /30, OUTPUT)
  DIMENSION X(5), PSI(30), TI(80), DF_4(30), DELV(30)
+, TEMF(10), W(200), IW(5)
   EXTERNAL LEG
   COMMON S. SETA, TH
   DELM(1)=0.0
   I=1
   M=0
15 PM=0.0
   PRINT*," "
   FRINT+," DELM=", TELM())
11 J=0
   L = 0
   T3=1200.0
   G=32.17405
   RE=2.092: 673E7
   GM=1.40734686E1E
   P1=3.1/15926536
10 Y(1) = 0.9
   X(2) = 0.0
   X(3) = .1
   X(4)=89.3*PI/180.0
   X(1) = 21030.0/6+SPRT(PM**2)+DELM(T)
   PM=6.04X(5)/7.0
   SM=X (5) -PM
   RETA=15.0/G
   TH=2.9"X(5)*G
   PRINT*," PARA",X(3),PM,SM,BETA,TH
   MEON=5
   ABS=1.05-8
   REL=1.0E-8
   EPS=1.0E-4
   JFLAG=1
   T=0.0
   IF (J.GT. 6) 60 TO 3
   TOUT=0.0
   DELT=TB/50.0
 1 TOUT=TOUT+SELT
   J=J+1
   CALL DOE(LEO, NERN, X, T, TOUT, REL, 175, IFLAG, 4, IW)
   JF(X(4).GT.0.3)GD TO 4
   IF(L.ED.1)GU TO !
   1 = 1
   DM=X(C)+SM
   IF (Dr.LT.0.0)50 TO 11
   DM=0.0
 . FSI(!)=X(4)
   TUOT=(U)IF
   7F(PCI(J).LT.-0.1) 30 TO 2
```

```
IF(TOUT.FO.T9) GO TO 5
    60 TO 1
 5 IF(PSI(J).LT.0.0130 TO 2
    GO TO 16
 2 K=4
   CALL INTERP(PSI,11, J, K, XO, YO, TEMP, IER)
   TF=Y()
   PRINT*," "
   PRINT","
              TIME=" .TE
    GO TO 10
 3 TOUT=TF
   CALL ODE (LED, NEON, Y, T, TOUT, REL, ABS, IFLAG, W, IN)
   PRINT*," X(I)",X(1),X(2),X(3),X(4),X(5),IOÚT
   DVOM=-TH/(RETA-X(3))+G/SIN(X(4))/BETA
   V=501 T (647 (RE+X(2)))
   PRINTY, "V CIRCULAR=".V
   TF (M.EC.1) GO TO 19
   PR1NT*," I=", [
   DELV(I)=X(3)-V
   JF(I.En.25)60 TO 18
   IF(I.GT.1)50 TO 17
   LEH=-DEFA(I) \UADA
   1=1+1
   DELM(I) = DEM+DELM(I-1)
   GO TC 15
17 I=I+1
   DELM(I)=DELM(I-1)+DEM
   GO TO 15
18 CONTINUE
   CALL INTERP(PELV, DELM, T, K, XO, YO, TEMP, IER)
   DELMF=YO
   PRINT*," "
PRINT*," DELTA MASS=", DELME
   T = 1
   v=1
   DELM(J)=DELME
   GO TO 15
16 PRINTE," PSI DOES NOT REACH OF
   STOP
   END
   SUCHOUTINE LEGIT, X. XP)
   DIMENSION X(_) .XP(3)
   COMMON G. 3FTA, 14
   XP(1)=X(3)+CCS(X(4))
   YP(2)=X(3)*SI4(Y(4))
   XP(3)=TH/X(5)-F1SIN(X(/))
   X^{p}(4) = -G^{c}COS(X(4))/X(3)
   YP(5)=-RETA
   RETUEN
  FNE
```

Appendix B

Impact Foint and Frobability Frediction Computer Program

```
PROGRAM FAIL (INPJT=/30,OUTPUT)
   DIMENSION X(5), X2(5), A7(12), EL(12), DV(3), W(205), IW(5)
  COMMON G, BETA, TH, 34, PI, RE
  REAL LAT, LATDA, LATDE, LATI, LONG, LONGO, LONGO, LONGT, LONGZ EXTERNAL LEQ
   K=0
   G=32.17405
  RE=2.0925673E7
  GM=1.40764688515
  PI=3.1415926536
  READ', TB, WE
   PRINT*,"
  PRINT*," T8=",T8," MASS=",WE
  BETA=15.0/6
   TH=2.8+WE+G
  PRINT+." BETA=". BETA."
                            THRUST= ". TH
  TOUT=TB
   IF (K.EQ. 0) GO TO 20
24 K=K+1
20 IFLAG=1
  X(1) = G.0
  X(2) = 0.0
  X(3) = .1
  X(4) = 89.8 + PI/180.0
  X(5) = WE
  NEQN=5
  AB=1.0E-8
  REL=AB
  T=0.0
  DELT=TB/100.0
  IF (K.EO.0) GO TO 21
   TOUT=TOUT-DELT
21 CALL ODE(LEG, NERN, X, T, TOUT, REL, A3, IFLAG, W, IW)
  PRINT+,"
              RANGE
                             ALT
                                            VEL
                                                          PSI
                                      TIME"
                 MA SS
  PRINT+, X(1), X(2), X(3), X(4), X(5), T) JT
  XP(1) = X(3) + COS(X(4))
  XP(2) = X(3) + SIV(X(1))
  PRINT+," XP1=",xP(1)," XP2=",XP(2)
  IF (K.GT.0) GO TO 25
  VCIRC=X(3)
  GO TC 24
25 CONTINUE
  DELV=VCIRC-X(3)
  FL (1) = -PI/2.0
  AZ(1)=0.0
```

```
DEL=FI/6.0
   DAZ=DEL
   DELA7=1.0E-5
   DELEL=1.0E-5
   DO 23 I=1,6
   DO 22 J=1,12
   L = 0
26 CONTINUE
   DV(1) = DEL V + COS (47(J)) + COS (EL (I))
   DV(2) = DEL V*SIN (A7(J)) * COS(EL(I))
   DV(3) = DELV+SIN(E_(I))
   IF(L.GT.0)GO TO 27
   PRINT+,"-----
   PRINT+," AZ=", AZ(J),"
                              EL=", EL(I)
27 CONTINUE
   CALL IMPACT(DV,X,XP,LAT,LONG,LONG?,DELT,L,ORBIT)
   IF(I.EQ.1)GO TO 32
   IF (A7(J).EQ.0.0) 30 TO 32
   EPS= . 001
   ZING= (AZ(J)-PI)*12
   IF (ZING.LT.EPS)G) TO 32
   IF(ORBIT.GT.0)GO TO 31
   IF(L.EQ.1) GO TO 23
   IF(L.EQ.2)GO TO 23
   L=1
   LATI=LAT
   LONGI = LONG 2
   AZI=FZ(J)
   ELI=FL(I)
   AZ(J) = AZI + DELAZ
   GO TO 26
28 CONTINUE
   LATOA=LAT
   LONGDA=LONG2
   DADLA=DELAZ/(LATDA-LATI)
   DADLO=DELAZ/(LONG)A-LONGI)
   L=2
   AZ(J) = AZI
   EL(I) = ELI + DELEL
   GO TO ?6
29 CONTINUE
   LAIDE=LAT
   LONG DE=LONG2
   DEDLA = DELEL/(LATDE-LATI)
   DEDLO=DELEL/(LONGOE-LONGI)
   EXP=ABS(COS(ELI))/(4.0'PI)
   P=ABS(DADLA*DEDLO-DADLO*DEDLA) *Ex2
   EL(I)=ELI
32 CONTINUE
   IF(L.GT.0)GO TO 33
   LATI=0.0
   P=0.0
33 CONTINUE
```

LATI=LATJ+180.0/°I

```
TIME=TOUT+DELT
    PRINT*,"
    PRINT+,"
                LATITUDE
                                LONGITUDE
                                              PROB43ILITY
                                                                 TIME"
    PRINT*, LATI, LONG, P, TIMF
    IF(I.EQ.1) GO TO 23
31 CONTINUE
22 AZ (J+1) = AZ (J) + DAZ
23 EL (I+1)=EL (I)+DE_
   IF(K.LT.5) GO TO 24
PRINT*," DONE"
    STOP
    END
   SUBROUTINE LEQ(T,X,XP)
   DIMENSION X(5),XP(3)
   COMMON G.BETA, TH, SM, PI
   XP(1) = X(3) + \cos(X(4))
   XP(2) = X(3) + SIN(X(4))
   XP(3)=TH/X(5)-G*SIV(X(4))
   XP(4) \approx -G*COS(X(4))/X(3)
   XP(5) = -BETA
   RETURN
   END
   SUBROUTINE IMPACT(DV, X, XP, LAT, LON3, _ ONG2; DELT, L, ORBIT)
   COMMON G, BETA, TH, 34, PI, RE
   DIMENSION V(3), R(3), H(3), EC(3), X(3), XP(5), DV(3)
   REAL MAGE, MAGR, MAGH, MAGV
   REAL LONG, LONG1, LONG2, LAT
   U=GM
   V(1) = DV(1) + XP(1)
   V(2)=DV(2)
   V(3) = DV(3) + XP(2)
   R(1) = X(1)
   R(2) = 0.0
   R(3) = X(2) + RE
    H=R* V
   H(1)=R(1)+V(2)-R(2)+V(1)
   H(2) = R(2) + V(3) - R(3) + V(2)
   H(3) = R(3) + V(1) - R(1) + V(3)
   H2 = 0
   R2=0
   V2=0
   COEFV=0
   00 11 I=1.3
   H2=H(I)**2*H2
   R2=R(I) ++ 2+R2
   V2=V(I)++2+V2
11 COEFV=R(I) +V(I)+COEFV
    GET ECCENTRICITY EC
   MAGH=SORT (H2)
   MAGR=SORT (R2)
   MAGV=SORT (V2)
   COEFF = MAGV + + 2-U/ 44 GR
```

```
E2=0
   DO 12 I=1.3
   EC(I)=(1.0/U) + (C)EFR+R(I) -COEFV+V(I))
12 E2=EC(I)+*2+E2
   MAGE=SORT (E2)
    TO FIND IF IT IMPACTS
   P=MAGH++2/U
   RPER=P/(1.0+M4GE)
   IF (RFER.GT.RE) GO TO 15
   A=P/(1.0-MAGE* *2)
   COEFTNU= ((P/MAGR) -1.0)/MAGE
   TANOM=ACOS (COEFT NJ)
   IF(CCEFV.GT.0.0) 30 TO 16
   TANOM=2.0*PI-TANOM
16 COEFENU= ((P/RE)-1.0) / MAGE
   EANOM=ACOS (COEFENJ)
   ANG=2.0*PI-TANOM-EANOM
   TANG=2.0* PI-EANO 4
   E=ACOS((MAGE+3OS(TANG))/(1.0+MAGE+3)S(TANG)))
   E0=ACOS((MAGE+COS(TANOM))/(1.0+MAGE+COS(TANOM)))
   IF (TANG.LT.PI) GO TO 17
   E=2.0*PI-E
17 IF (TANOM.LT.PI)30 TO 18
   E0=2.0*PI-E0
18 CONTINUE
   DELT=SQRT(Af+3/U)+((E-MAGE+SIN(E))-(E0-MAGE+SIN(E0)))
   AINC=ACOS (H(3) /MAGH)
   B CT CD(0.0.TD.(2) TO 8
   AINC = - AINC
 8 CONTINUE
   LONG1=ASIN(X(1)/RE)
   A180=PI
   IF (AINC.EQ.0.0)33 TO 9
   IF (AINC. EQ. A 18 0) 30 TO 9
   LAT=ASIN(SIN(AINC) *SIN(ANG))
   LONG Z=ASIN (TAN (LAT) / TAN (AINC))
   CALL REDUCE(ANG, LONG?)
   GO TO 19
 9 CONTINUE
   LONG 2=ANG
   LAT=0.0
19 CONTINUE
   LONG= (LONG1+LONG2) *180.0/PI
   ANG=ANG+180.0/PI
   AINC=AINC+180.0/PI
   IF(L.GT.0)GO TO 10
PRINT*," LAT=",LAT,"LONG=",LONG2
   PRINT*,"
   PRINT*,"
                                TIME
                                                INCLINATION"
                ANGLE
   PRINT*, ANG, DELT, AINC
10 CONTINUE
   ORBIT=0.0
   GO TC 14
15 PRINT*,"
             ORBIT MAS ECCENTRICITY=", MAGE
```

ORBIT=1.0 14 RETUPN END SUBROUTINE REDUCE(ANS, LONG2) REAL LONG 2 PI=3.1415926535 A90=PI/2.0 A180=PI A270=1.5*PI A360=2.0*PI IF (ANG.LT.A360) 33 TO 8 ANI=ANG/A360 IAN=INT(ANI) ANG= (ANI-IAN) + 2.0 + PI 8 CONTINUE IF (ANG.LT.A90) GO TO 7 IF (ANG.LT.A180)50 TO 6 IF (ANG.LT.A270)50 TO 6 LONG 2=A360+LONG2

GO TO 4

- 6 LONG 2=A180-LONG2
- 4 RETURN END

<u>Vita</u>

Scott W. Benson was born on 8 November 1955 in Geneva, Illinois. He graduated from Geneva Community High School in 1973. In May 1977, he received the degree of Bachelor of Science in Aeronautical and Astronautical Engineering from Purdue University. He received his commission through Air Force ROTC in May 1977. He entered the Air Force Institute of Technology in September 1977 as his first active duty assignment. Upon completion of his studies, he was assigned to the Foreign Technology Division, FTD/SDSY, at Wright-Patterson Air Force Base.

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SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

REPORT DOCUMENTATION I	READ INSTRUCTIONS BEFORE COMPLETING FORM				
1. REPORT NUMBER AFIT/GA/AA/78D-1	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER			
4. TITLE (and Subtitle)	5. TYPE OF REPORT & PERIOD COVERED				
PROBABILITY DISTRIBUTION OF THE FOR A SATELLITE BOOSTER AFTER	MS Thesis				
OF CONTROL	6. PERFORMING ORG. REPORT NUMBER				
Scott W. Benson, B.S. 2nd Lt USAF		8. CONTRACT OR GRANT NUMBER(s)			
9. PERFORMING ORGANIZATION NAME AND ADDRESS Air Force Institute of Technology AFTT/EN Wright-Patterson AFR, Ohio Ly	10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS				
11. CONTROLLING OFFICE NAME AND ADDRESS		12. REPORT DATE March 1979 19. NUMBER OF PAGES 40			
14. MONITORING AGENCY NAME & ADDRESS(If different	15. SECURITY CLASS. (of this report) Unclassified 15a. DECLASSIFICATION DOWNGRADING SCHEDULE				
6. DISTRIBUTION STATEMENT (of this Report)					

Approved for public release; distribution unlimited

17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)

18. SUPPLEMENTARY NOTES

Approved for public release, IAW APR 190-17

JOSEFH F. HIRFS, Major, USAF Director of Information

19. KEY WORDS (Continue on reverse side if necessary and identify by block number)

Booster Failure Impact Frediction Impact Frobability Scacecraft Launch

20. ABSTRACT (Continue on reverse side if necessary and identity by block number)

The impact area and regions of high impact probability after booster failure were found for a sample space system. The impact footprints were found to be drop-shared. They increased in size as the residual velocity vector elevation approached the original flight path, or when the booster failure was assumed earlier in the boost trajectory. Probability densities at impact were found, and plotted against the longitude of impact. Singularities in the probability density were discovered. The singularities correspond

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Unclassified SECURITY CLASSIFICATION OF THIS PAGE (When Date Entered) to areas of high impact probability, and occur when a maximum latitude or longitude is obtained at a fixed residual velocity vector azimuth or elevation. These singularities were plotted along the impact envelopes. A sample failing booster system was used to demonstrate the impact area and probability characteristics.

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